

PHOTONICS Research

Deterministic distribution of orbital angular momentum multiplexed continuous-variable entanglement and quantum steering

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Orbital angular momentum (OAM) multiplexing provides an efficient method to improve data-carrying capacity in various quantum communication protocols. It is a precondition to distribute OAM multiplexed quantum resources in quantum channels for implementing quantum communication. However, quantum steering of OAM multiplexed optical fields and the effect of channel noise on OAM multiplexed quantum resources remain unclear. Here, we generate OAM multiplexed continuous-variable (CV) entangled states and distribute them in lossy or noisy channels. We show that the decoherence property of entanglement and quantum steering of the OAM multiplexed states carrying topological charges $l = 1$ and $l = 2$ are the same as that of the Gaussian mode with $l = 0$ in lossy and noisy channels. The sudden death of entanglement and quantum steering of high-order OAM multiplexed states is observed in the presence of excess noise. Our results demonstrate the feasibility to realize high data-carrying capacity quantum information processing by utilizing OAM multiplexed CV entangled states. © 2022 Chinese Laser Press

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1. INTRODUCTION

Einstein–Podolsky–Rosen (EPR) entanglement plays a crucial role in quantum information processing, such as quantum communication, quantum computation, and quantum precision measurement [1–5]. Besides entanglement, quantum steering, which stands between entanglement [1] and Bell nonlocality [6] in the hierarchy of quantum correlations [7], has been identified as a useful quantum resource. Different from entanglement and Bell nonlocality, quantum steering shows unique asymmetry or even one-way characteristics [8–16] and, thus, allows asymmetric quantum information processing. For example, quantum steering enables one-side device-independent quantum key distribution [17–19].

Multiplexing provides an efficient method to enhance the data-carrying capability in both classical and quantum communication systems by combining multiple channels into a single channel. By utilizing different degrees of freedom (DOFs) of light, such as wavelength [20,21], polarization [22], temporal [23–25] or spatial [26,27] modes, different types of multiplexing can be realized. Orbital angular momentum

(OAM) of light [28] has also been found to be an attractive DOF to realize multiplexing due to its infinite range of possibly achievable topological charges [29,30]. OAM has found applications in discrete-variable quantum information processing, such as high-dimensional OAM entanglement generation [31], and 18-qubit entanglement with six photons' three DOFs including OAM [32].

Four-wave mixing (FWM) process in warm alkali vapor cell has found a wide range of applications [33–39]. Especially, spatial-multi-mode advantage of the FWM process, attributed to its cavity-free configuration, makes it an ideal optical parametric amplifier to generate entangled images [35] and reconfigurable multipartite entanglement [36]. Quantum correlated twin beams carrying OAM were generated based on the FWM process in rubidium vapor [37]. OAM multiplexed bipartite and multipartite continuous-variable (CV) entangled states have been generated based on the FWM process [40–42]. Furthermore, OAM multiplexed deterministic all-optical quantum teleportation has also been demonstrated by utilizing the OAM multiplexed bipartite CV entangled state generated from

the FWM process [43]. To enhance the data-carrying capacity in quantum communication based on OAM multiplexed CV entangled states, it is essential to distribute them in lossy and noisy quantum channels towards practical applications. The distributions of weak coherent field and single photons carrying OAM in fiber, free space, and underwater have been experimentally investigated [44–46]. However, it remains unclear whether the quantum entanglement and steering of OAM multiplexed CV entangled states are more sensitive to loss and noise than a commonly used Gaussian mode with $l = 0$.

Here, we present the deterministic distribution of OAM multiplexed CV quantum entanglement and steering in lossy and noisy channels. In the experiment, the OAM multiplexed entangled fields are generated deterministically based on the FWM process in warm cesium vapor and distributed deterministically in quantum channels. We show that the CV entangled states carrying topological charges $l = 1$ and $l = 2$ are as robust against loss as the Gaussian mode with $l = 0$. Sudden death of entanglement and quantum steering of a high-order OAM multiplexed CV entangled state is observed in the presence of noise. Our results pave the way for applying OAM multiplexed CV entanglement and quantum steering in high data-carrying capacity quantum communication.

2. PRINCIPLE AND EXPERIMENTAL SETUP

Figure 1(a) shows the schematic of the experimental setup, and Fig. 1(b) shows the double- Λ energy-level structure used for the FWM process, which is formed from the ^{133}Cs D₁ line with an excited level ($6P_{1/2}, F' = 4$) and two ground levels

($6S_{1/2}, F = 3$ and $F = 4$). The pump beam is about 1.6-GHz blue detuned from $6S_{1/2}, F = 3 \rightarrow 6P_{1/2}, F' = 4$ transition, and the probe beam is 9.2-GHz redshifted relative to the pump beam. The pump and probe beams are combined by a GL and then cross each other in the center of the cesium vapor cell at an angle of 6 mrad [47]. The gain of the FWM process is around 3 with a pump power of 240 mW and a probe power of 3 μW . By injecting the probe beam carrying topological charge l of the OAM mode, a conjugate beam carrying topological charge $-l$ of the OAM mode is generated on the other side of the pump, which satisfies OAM conservation in the FWM process. The topological charge of the OAM mode $l = 1$ or $l = 2$ is added to the probe beam by passing it through a VPP. The pump beam is filtered out by using a GT with an extinction ratio of $10^5:1$ after the vapor cell.

The Conj field is kept by Alice, whereas the Pr field is distributed to a remote quantum node owned by Bob through a lossy or noisy channel. The lossy channel is simulated by an HWP and a PBS. The noisy channel is modeled by combining the Pr field with an auxiliary beam at a PBS followed by an HWP and a PBS. The auxiliary beam carries the same frequency and topological charge as the Pr field and is modulated by an AM and a PM with white noise [16]. To characterize the OAM multiplexed CV entangled state, its covariance matrix is experimentally measured by utilizing two sets of BHDs. In order to extract the CV quadrature information carried by the OAM mode with a topological charge l , an LO with opposite topological charge $-l$ is required. In our experiment, the spatially mode-matched LO beams used in the BHDs are obtained from a second set of FWM processes which is around 5 mm

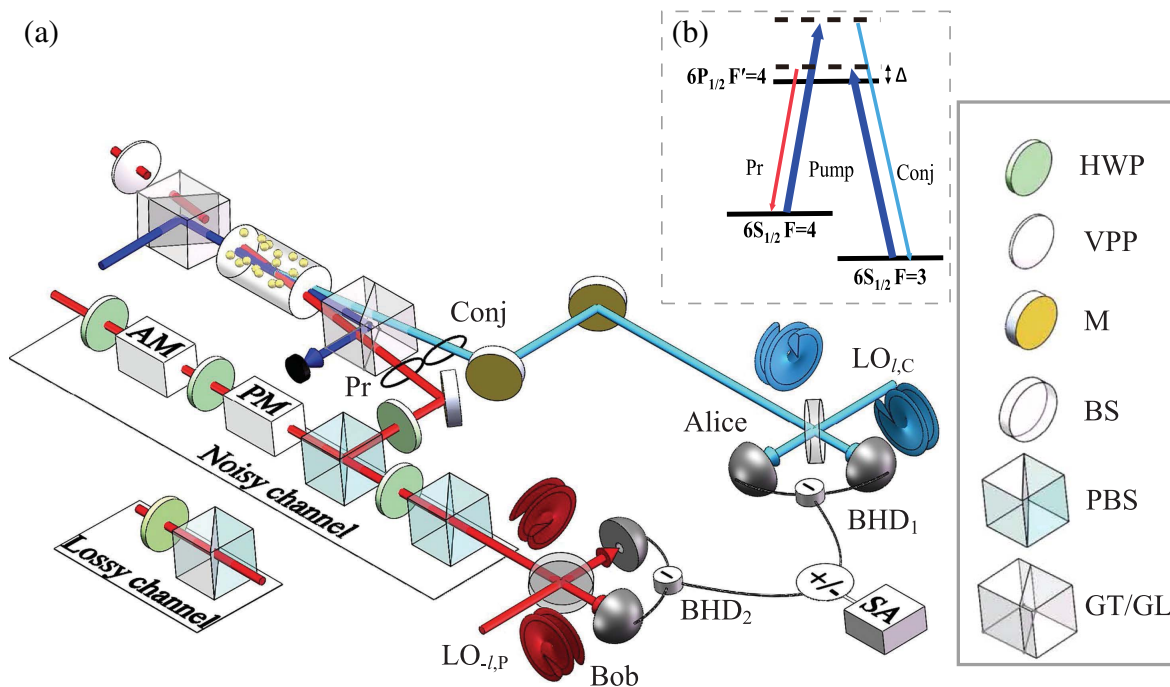


Fig. 1. (a) Experimental setup for the generation and distribution of OAM multiplexed CV quantum entanglement and steering in a lossy or noisy channel. Pr, probe beam; Conj, conjugate beam; $\text{LO}_{-l,P}$ and $\text{LO}_{l,C}$, local oscillators of Pr and Conj fields; AM, amplitude modulator; PM, phase modulator; GL, Glan-laser polarizer; GT, Glan-Thompson polarizer; PBS, polarization beam splitter; HWP, half-wave plate; VPP, vortex phase plate; M, mirror; BS, 50:50 beam splitter; BHD₁, BHD₂, balanced homodyne detectors; SA, spectrum analyzer. (b) Double- Λ energy-level structure for the FWM process in a cesium vapor cell. Δ , one-photon detuning.

above the first set of FWM processes in the same vapor cell [35]. More details of the experimental parameters can be found in Appendix B.

The Hamiltonian of the OAM multiplexed FWM process can be expressed as [40]

$$\hat{H} = \sum_l i\hbar\gamma_l \hat{a}_{l,p}^\dagger \hat{a}_{-l,C}^\dagger + \text{H.c.}, \quad (1)$$

where γ_l , which is defined as the interaction strength of each OAM pair, is proportional to the pump power. $\hat{a}_{l,p}^\dagger$ and $\hat{a}_{-l,C}^\dagger$ are the creation operators related to OAM modes of the Pr and Conj fields, respectively, and H.c. represents the Hermitian conjugate. Since the pump beam does not carry OAM ($l = 0$), the topological charges of the Pr and Conj fields are opposite due to OAM conservation in the FWM process. The output state of the OAM multiplexed FWM process is as follows:

$$|\Psi\rangle_{\text{out}} = |\Psi\rangle_{-l} \otimes \cdots \otimes |\Psi\rangle_0 \otimes \cdots \otimes |\Psi\rangle_l, \quad (2)$$

where $|\Psi\rangle_l = |\psi_{l,p}, \psi_{-l,C}\rangle$ presents a series of independent OAM multiplexed CV EPR entangled states generated in the FWM process. $|\psi_{l,p}\rangle$ and $|\psi_{-l,C}\rangle$ represent Pr field carrying topological charge l and Conj field carrying topological charge $-l$, respectively.

All Gaussian properties of the CV entangled state $|\Psi\rangle_l$ can be determined by the covariance matrix σ_{AB} with matrix element $\sigma_{ij} = \langle \hat{\xi}_i \hat{\xi}_j + \hat{\xi}_j \hat{\xi}_i \rangle / 2 - \langle \hat{\xi}_i \rangle \langle \hat{\xi}_j \rangle$, where $\hat{\xi} \equiv (\hat{X}_{-l,C}, \hat{Y}_{-l,C}, \hat{X}_{l,p}, \hat{Y}_{l,p})^T$, $\hat{X} = \hat{a} + \hat{a}^\dagger$, and $\hat{Y} = (\hat{a} - \hat{a}^\dagger)/i$ represent amplitude and phase quadratures of the Conj and Pr fields, respectively. The covariance matrix of the OAM multiplexed entangled state after distribution in a lossy or noisy channel is as follows:

$$\sigma_{AB,\delta,\eta} = \begin{pmatrix} V_a & 0 & V_c & 0 \\ 0 & V_a & 0 & -V_c \\ V_c & 0 & V_b & 0 \\ 0 & -V_c & 0 & V_b \end{pmatrix}, \quad (3)$$

with $V_a = \frac{V+V'}{2}$, $V_b = \eta \frac{V+V'}{2} + (1-\eta)(1+\delta)$, and $V_c = \sqrt{\eta} \frac{V'-V}{2}$. V and V' represent the variances of squeezed and antisqueezed quadratures of the optical mode, respectively. The derivation of the covariance matrix can be found in Appendix A. For pure squeezed states, $VV' = 1$, whereas $VV' > 1$ after pure squeezed states suffer from loss or noise. η and δ represent transmission efficiency and excess noise of the quantum channel, respectively. We have $\delta = 0$ in a lossy channel, whereas we have $\delta > 0$ in a noisy channel. The submatrices $\sigma_A = V_a I$ and $\sigma_B = V_b I$ with I as the identity matrix are the covariance matrices corresponding to the states of Alice's and Bob's subsystems, respectively. The CV entangled state is a symmetric state and an asymmetric state when $\sigma_A = \sigma_B$ and $\sigma_A \neq \sigma_B$, respectively.

The Peres–Horodecki criterion of positivity under the partial transpose (PPT) criterion is a sufficient and necessary criterion to characterize the entanglement of CV bipartite entanglement [48]. If the smallest symplectic eigenvalue ν of the partially transposed covariance matrix is smaller than 1, bipartite entanglement exists. Otherwise, it is a separable state. Furthermore, smaller ν represents stronger entanglement.

Quantum steering for bipartite Gaussian states of CV systems can be quantified by [49]

$$\mathcal{G}^{A \rightarrow B}(\sigma_{AB}) = \max \left\{ 0, \frac{1}{2} \ln \frac{\det \sigma_A}{\det \sigma_{AB}} \right\}, \quad (4)$$

where $\mathcal{G}^{A \rightarrow B}(\sigma_{AB}) > 0$ represents that Alice has the ability to steer Bob's state. Similarly, we have

$$\mathcal{G}^{B \rightarrow A}(\sigma_{AB}) = \max \left\{ 0, \frac{1}{2} \ln \frac{\det \sigma_B}{\det \sigma_{AB}} \right\}, \quad (5)$$

which represents Bob's ability to steer Alice's state. From the expressions $\mathcal{G}^{A \rightarrow B}(\sigma_{AB})$ and $\mathcal{G}^{B \rightarrow A}(\sigma_{AB})$, it can be seen that Alice and Bob have the same steerability if $\det \sigma_A = \det \sigma_B$ is satisfied, i.e., the bipartite Gaussian state is a symmetric state. If the state is an asymmetric state, the steerabilities of Alice and Bob will be different.

3. RESULTS

To verify the OAM property of the optical fields, we measure the spatial beam patterns of quantum states $|\Psi\rangle_1$ and $|\Psi\rangle_2$ transmitted through a lossy channel, which are shown in the top rows of Figs. 2(a) and 2(b), respectively. It is obvious that the Pr and Conj fields are both Laguerre–Gaussian beams. To infer their topological charges, they are passed through a tilted lens and imaged on a camera. As shown in the bottom rows of Figs. 2(a) and 2(b), the number of dark stripes gives the number of the topological charge, and the direction gives its sign [50]. As the transmission efficiency of the Pr field decreases, its optical intensity also decreases, whereas its topological charge remains unchanged. Additional beam patterns of the Pr and Conj fields can be found in Appendix D.

The covariance matrices of the OAM multiplexed entangled states are reconstructed by measuring the noise variances of the amplitude and phase quadratures of the Conj and Pr fields $\Delta^2 \hat{X}_{-l,C}$, $\Delta^2 \hat{Y}_{-l,C}$, $\Delta^2 \hat{X}_{l,p}$, $\Delta^2 \hat{Y}_{l,p}$, as well as their correlation variances of amplitude and phase quadratures $\Delta^2(\hat{X}_{l,p} - \hat{X}_{-l,C})$ and $\Delta^2(\hat{P}_{l,p} + \hat{P}_{-l,C})$, respectively. Details about the measurement of the covariance matrices can be found in Appendix C. Based on the covariance matrix of each OAM multiplexed entangled state at different loss and noise levels, its quantum entanglement and quantum steering characteristics are evaluated experimentally.

Figure 2(c) shows the dependence of PPT values of the CV bipartite entangled state carrying different topological charges on the transmission efficiency of the Pr field. The correlation and anticorrelation levels of the initial CV entangled states carrying topological charges $l = 0$, $l = 1$, and $l = 2$ are all around -3.3 and 6.1 dB, which correspond to $V = 0.47$ and $V' = 4.11$, respectively. The entanglement between Pr and Conj fields degrades as the transmission efficiency decreases. However, the entanglement is robust against loss, i.e., it always exists until the transmission efficiency reaches 0. It is obvious that the CV bipartite entangled state carrying topological charges $l = 1$ and $l = 2$ is as robust to loss as its Gaussian counterpart $l = 0$.

Figure 3 shows the dependence of PPT values of the CV bipartite entangled state carrying different topological charges in noisy channels. Compared with the results in Fig. 2(c), the

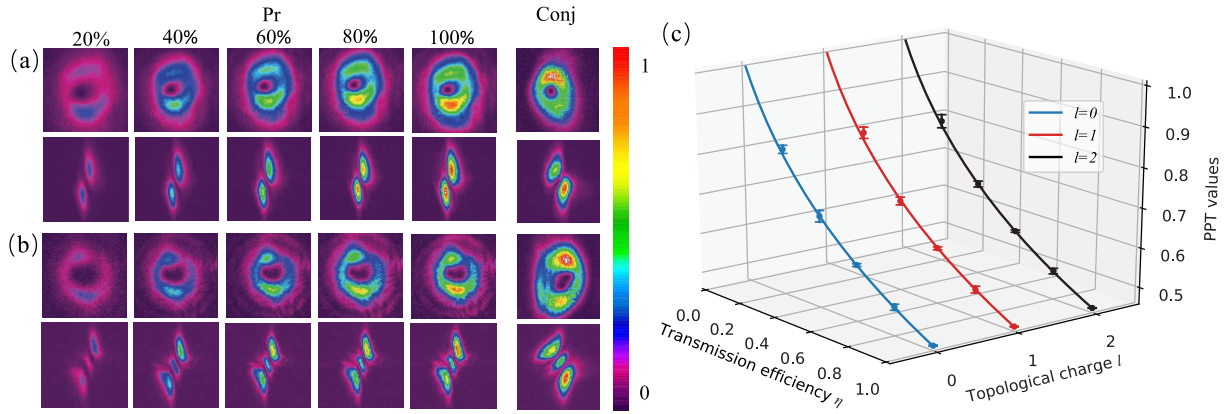


Fig. 2. (a) and (b) Beam patterns of the OAM multiplexed CV entanglement for $l = 1$ and $l = 2$ in a lossy channel and corresponding transmitted patterns through a tilted lens. (c) Dependence of PPT values of the OAM multiplexed CV entanglement on transmission efficiency η for $l = 0$, $l = 1$, and $l = 2$ in lossy channels. The PPT value is 0.46 ± 0.01 at $\eta = 1$. Curves and data points show theoretical predictions and experimental results, respectively. Error bars of experimental data represent one standard deviation and are obtained based on the statistics of the measured data.

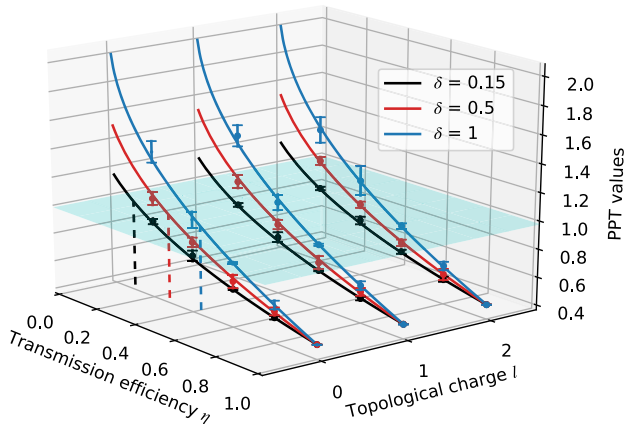


Fig. 3. Dependence of PPT values of the OAM multiplexed CV entanglement on transmission efficiency η for $l = 0$, $l = 1$, and $l = 2$ in noisy channels. Three different amounts of excess noise $\delta = 0.15$ (black), $\delta = 0.5$ (red), and $\delta = 1$ (blue) are compared. The light blue plane shows the boundary for sudden death of entanglement where the PPT value equals 1. The three vertical dashed lines indicate corresponding transmission efficiencies where entanglement starts to disappear. Curves and data points show theoretical predictions and experimental results, respectively. Error bars of experimental data represent one standard deviation and are obtained based on the statistics of the measured data.

entanglement disappears at a certain transmission efficiency of the Pr field in the presence of excess noise, which demonstrates the sudden death of CV quantum entanglement. Furthermore, the higher the excess noise is, the sooner entanglement disappears. The transmission efficiencies where entanglement starts to disappear are $\eta = 0.10$, 0.28 , and 0.44 , respectively, for the excess noise $\delta = 0.15$, 0.5 , and 1 in the units of shot-noise level (SNL). We show that OAM multiplexed CV entangled states carrying high-order topological charges $l = 1$ and $l = 2$ exhibit the same decoherence tendencies as their Gaussian counterpart $l = 0$ in noisy channels.

The dependence of steerabilities $\mathcal{G}^{A \rightarrow B}$ and $\mathcal{G}^{B \rightarrow A}$ on the transmission efficiency η and topological charge l in lossy and noisy channels is shown in Figs. 4(a) and 4(b), respectively. In a lossy channel, the steerabilities for both directions always decrease when the transmission efficiency decreases. One-way steering is observed in the region of $0 < \eta < 0.72$ for the OAM multiplexed CV entangled state carrying different topological charges $l = 0$, $l = 1$, and $l = 2$. In a noisy channel where the excess noise $\delta = 0.15$ (in the units of SNL) exists, the steerabilities $\mathcal{G}^{A \rightarrow B}$ and $\mathcal{G}^{B \rightarrow A}$ are lower than those in the lossy channel [13,51]. Furthermore, Alice loses steerability in the region of $0 < \eta < 0.45$, whereas Bob loses steerability in the region of $0 < \eta < 0.75$, which confirms sudden death of quantum steering in a noisy channel [51]. It is worth noting that the CV entangled state carrying topological charges $l = 1$ and $l = 2$ has the same steerabilities as its counterpart $l = 0$.

4. CONCLUSION

The distribution of OAM multiplexed CV entanglement and quantum steering in quantum channels with homogeneous loss and noise, such as fiber channels, was experimentally simulated in our paper. There were also other quantum channels with inhomogeneous loss and noise, such as atmospheric turbulence and diffraction. Recently, it was shown that other optical fields carrying OAM, such as vector beams, were turbulence resilient in atmospheric turbulence [52]. Thus, it is worthwhile to investigate the turbulence-resilient characteristics of OAM multiplexed CV quantum entanglement and steering, which have the potential to substantially improve the quantum communication distance and fidelity.

To summarize, we experimentally demonstrated quantum steering of OAM multiplexed optical fields and investigated the distribution of OAM multiplexed CV entanglement and quantum steering in quantum channels. We showed that the decoherence property of CV entanglement and quantum steering of the OAM multiplexed optical fields carrying topological charges $l = 1$ and $l = 2$ were the same as that of the

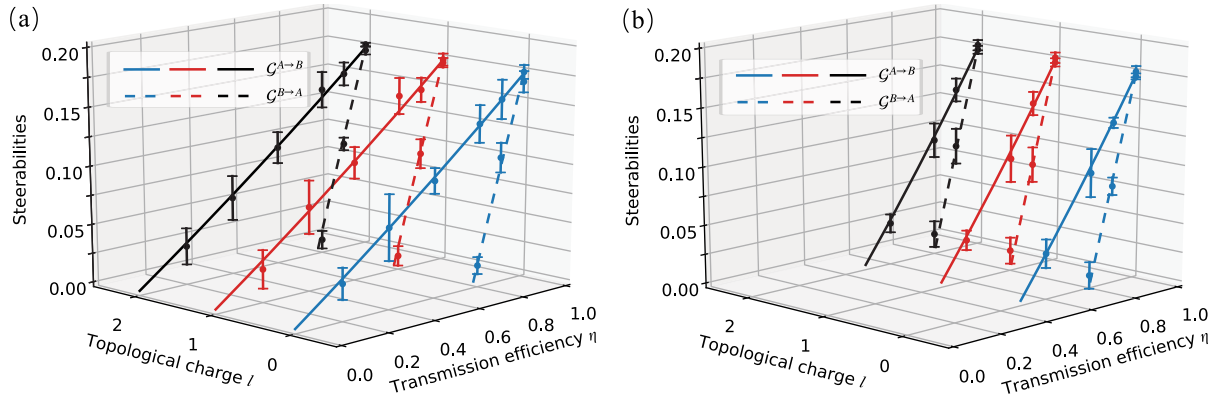


Fig. 4. Quantum steerabilities of OAM multiplexed CV entangled state distributed in a (a) lossy or (b) noisy channel. The excess noise shown in (b) is $\delta = 0.15$. Solid and dashed curves show theoretical predictions of $G^{A \rightarrow B}$ and $G^{B \rightarrow A}$, respectively. Data points show experimental results. Error bars of experimental data represent one standard deviation and are obtained based on the statistics of the measured data.

counterpart Gaussian mode with $l = 0$ in lossy and noisy channels. The sudden death of entanglement and quantum steering of high-order OAM multiplexed optical fields was observed in the presence of excess noise. Our results demonstrated the feasibility to improve the quantum communication capacity in practical quantum channels by utilizing OAM multiplexed CV entanglement and quantum steering.

APPENDIX A: THEORETICAL MODEL

The Hamiltonian of the OAM multiplexed FWM process can be expressed as

$$\hat{H} = \sum_l i\hbar k_l \hat{a}_{0,\text{Pump}} \hat{a}_{0,\text{Pump}} \hat{a}_{l,P}^\dagger \hat{a}_{-l,C}^\dagger e^{i\theta} + \text{H.c.}, \quad (\text{A1})$$

where $\hat{a}_{0,\text{Pump}}$, $\hat{a}_{l,P}^\dagger$, and $\hat{a}_{-l,C}^\dagger$ are the annihilation operator of the pump field and the creation operators related to OAM modes of the Pr and Conj fields, respectively. θ is the phase of the pump field. The first term represents the process in which two pump photons are converted into a Pr photon and a Conj photon. k_l is the interaction strength of each OAM pair. Since the pump beam does not carry OAM ($l = 0$), the topological charges of the Pr and Conj fields are opposite due to OAM conservation in the FWM process. The second term H.c. is the H.c. of the first term and represents the reversed process in which a Pr photon and a Conj photon are converted to two pump photons.

The pump field is much stronger than the Pr and Conj fields in the FWM process, so it can be regarded as a classical field. By combining the intensity of the pump field $|\alpha_{\text{pump}}|^2$ with k_l , i.e., $|\alpha_{\text{pump}}|^2 k_l = \gamma_l$, and taking $\theta = 0$, the Hamiltonian can be simplified as [40]

$$\hat{H} = \sum_l i\hbar \gamma_l \hat{a}_{l,P}^\dagger \hat{a}_{-l,C}^\dagger + \text{H.c.} \quad (\text{A2})$$

The output state of the OAM multiplexed FWM process is as follows:

$$|\Psi\rangle_{\text{out}} = |\Psi\rangle_{-l} \otimes \cdots \otimes |\Psi\rangle_0 \otimes \cdots \otimes |\Psi\rangle_l, \quad (\text{A3})$$

where $|\Psi\rangle_l = \hat{S}(\gamma_l)|\text{vac}\rangle = |\psi_{l,P}, \psi_{-l,C}\rangle$ means the EPR entangled state with Pr and Conj fields carrying topological

charges l and $-l$, respectively. $|\text{vac}\rangle$ represents the vacuum state, and $\hat{S}(\gamma_l) = e^{\gamma_l(\hat{a}_{l,P}^\dagger \hat{a}_{-l,C}^\dagger - \hat{a}_{l,P} \hat{a}_{-l,C})}$ represents the two-mode squeezing operator for the vacuum state carrying OAM topological charge l . It is obvious that a series of CV EPR entangled states carrying independent topological charges is generated in the FWM process; i.e., OAM multiplexing is realized.

All Gaussian properties of the CV Gaussian entangled state $|\Psi\rangle_l$ can be determined by its covariance matrix σ_{AB} with the matrix element $\sigma_{ij} = \langle \hat{\xi}_i \hat{\xi}_j + \hat{\xi}_j \hat{\xi}_i \rangle / 2 - \langle \hat{\xi}_i \rangle \langle \hat{\xi}_j \rangle$, where $\hat{\xi} \equiv (\hat{X}_{-l,C}, \hat{Y}_{-l,C}, \hat{X}_{l,P}, \hat{Y}_{l,P})^T$, and $\hat{X} = \hat{a} + \hat{a}^\dagger$ and $\hat{Y} = (\hat{a} - \hat{a}^\dagger)/i$ represent amplitude and phase quadratures of the Conj and Pr fields.

The covariance matrix of the CV bipartite entangled state can be written as

$$\begin{aligned} \sigma_{AB} &= \begin{pmatrix} \frac{V+V'}{2} & 0 & \frac{V'-V}{2} & 0 \\ 0 & \frac{V+V'}{2} & 0 & \frac{V-V'}{2} \\ \frac{V'-V}{2} & 0 & \frac{V+V'}{2} & 0 \\ 0 & \frac{V-V'}{2} & 0 & \frac{V+V'}{2} \end{pmatrix} \\ &= \begin{pmatrix} \frac{V+V'}{2} I & \frac{V'-V}{2} Z \\ \frac{V'-V}{2} Z & \frac{V+V'}{2} I \end{pmatrix}, \end{aligned} \quad (\text{A4})$$

where V and V' represent the variances of squeezed and anti-squeezed quadratures of the optical mode, respectively. I and Z are the Pauli matrices,

$$I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \quad (\text{A5})$$

Then, we consider the distribution of CV entangled state $|\Psi\rangle_l$ in a lossy or noisy channel. Let $\hat{a}_{l,P}$ and $\hat{a}_{-l,C}$ represent the annihilation operators of the Pr and Conj fields, respectively. After the Pr field $\hat{a}_{l,P}$ is distributed in a lossy channel, it becomes $\hat{a}'_{l,P} = \sqrt{\eta} \hat{a}_{l,P} + \sqrt{1-\eta} \hat{\mu}$, where $\hat{\mu}$ represents the vacuum state with a variance of 1. Similarly, the Pr field becomes $\hat{a}'_{l,P} = \sqrt{\eta} \hat{a}_{l,P} + \sqrt{1-\eta} (\hat{e} + \hat{\mu})$ after it is distributed in a noisy channel with excess noise $\Delta^2(\hat{X}_e) = \Delta^2(\hat{Y}_e) = \delta$ [13]. $\delta = 0$ means that there is no excess noise, and only loss exists in the channel. $\delta > 0$ means that there exists excess noise

in the channel. So the covariance matrix of CV entangled state $|\Psi\rangle_l$ after distribution in the lossy or noisy channel is as follows:

$$\sigma_{AB,\delta,\eta} = \begin{pmatrix} V_a & 0 & V_c & 0 \\ 0 & V_a & 0 & -V_c \\ V_c & 0 & V_b & 0 \\ 0 & -V_c & 0 & V_b \end{pmatrix} = \begin{pmatrix} V_a I & V_c Z \\ V_c Z & V_b I \end{pmatrix}, \quad (\text{A6})$$

with $V_a = \frac{V+V'}{2}$, $V_b = \eta \frac{V+V'}{2} + (1-\eta)(1+\delta)$, and $V_c = \sqrt{\eta} \frac{V'-V}{2}$.

APPENDIX B: DETAILS OF THE EXPERIMENT

The Ti:sapphire laser (Coherent MBR-110) is about 1.6 GHz blue detuned from the ^{133}Cs D1 line $6S_{1/2}, F = 3 \rightarrow 6P_{1/2}, F' = 4$ transition with a total power of 1.2 W. As shown in Fig. 5, the laser beam is split into two beams by PBS₁. The horizontally polarized beam is split into two beams by PBS₂, and they serve as the pump beams for two sets of FWM processes in the same cesium vapor cell. The vertically polarized beam with a power of 30 mW passes through a resonance EOM (Qubig GmbH PM-Cs) and three successive temperature-stabilized etalons to realize a 9.2-GHz frequency redshift [47]. The OAM of topological charge $l = 1$ or $l = 2$ is added to the frequency redshifted beam by passing it through a VPP (RPC Photonics). Then the frequency redshifted beam is split into three beams, among which the first two beams serve as the probe beams of two FWM processes, and the third beam serves as the auxiliary beam for the noisy channel. The pump and probe beams are combined in a GL and then cross each other in the center of the cesium vapor cell (Traid

Technology, Inc.) at an angle of 6 mrad. The vapor cell is 25-mm long, and its temperature is stabilized at 103°C.

The two sets of FWM processes are constructed in the same cesium vapor cell with a height difference of 5 mm. The bottom FWM process is used to generate the OAM multiplexed CV entangled state, whereas the top FWM process is used to generate spatially matched LOs with Pr and Conj fields. The pump power in the bottom FWM process for generating OAM multiplexed CV entanglement is 240 mW. The probe gain of the bottom FWM process is around 3, and the degree of initially generated CV entanglement is around -3.3 ± 0.1 dB. The pump power and seed probe power of the top FWM processes are 450 mW and 100 μW , respectively, so that the SNL is around 10 dB higher than the electronic noise of the homodyne detector. The bottom FWM process is weakly seeded with a probe power of around 3 μW for relative phase locking of the Pr/Conj fields and their LOs in the balanced homodyne detections.

The lossy channel is simulated by HWP₁₀ and PBS₆. The noisy channel is modeled by combining the vertically polarized Pr field with a horizontally polarized auxiliary beam at PBS₅ followed by HWP₁₀ and PBS₆. The auxiliary beam carries the same frequency and topological charge as the Pr field and is modulated by an AM and a PM with excess noise at 1.2 MHz. The amount of excess noise is adjusted by tuning the amplitude of the signal applied to the AM and PM and evaluated in the units of SNL. For example, excess noise $\delta = 1$ corresponds to a noise level that is 3 dB higher than the SNL. By tuning HWP₁₀, the lower the transmission efficiency of the Pr field, the higher the excess noise coupled to the Pr field. In practical quantum communication protocols, higher excess noise is coupled to the quantum entangled state

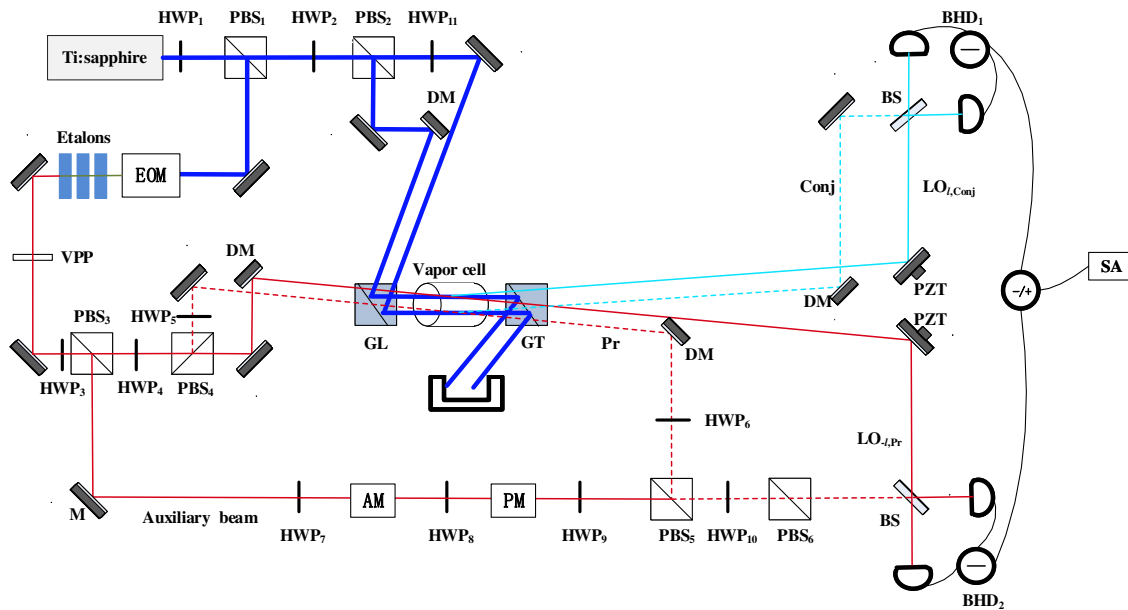


Fig. 5. Detailed experimental schematic for distributing OAM multiplexed CV entanglement in a noisy channel. The lossy channel is realized by blocking the auxiliary beam. D-shaped mirrors (DMs) are utilized to combine or separate light beams with small distances. HWP, half-wave plate; PBS, polarization beam splitter; EOM, electro-optic modulator; VPP, vortex phase plate; GL, Glan-laser polarizer; GT, Glan-Thompson polarizer; Pr, probe beam; Conj, conjugate beam; AM, amplitude modulator; PM, phase modulator; M, mirror; DM, D-shaped mirror; PZT, piezoelectric ceramics; BS, 50:50 beam splitter; BHD, balanced homodyne detector; SA, spectrum analyzer.

as the communication distance increases, accompanying lower transmission efficiency. Therefore, our experimental setting is similar to the realistic scenarios in the practical noisy quantum channel. To characterize the OAM multiplexed CV entangled state, its covariance matrix is experimentally obtained by utilizing two sets of BHDs (Thorlabs PDB450A). The interference visibilities for the two sets of BHDs are both around 99%. The electrical gains of these two BHDs are both 10^5 V/A. The original photodiodes are replaced by high quantum efficiency (QE) photodiodes with QE = 98% at 895 nm (first sensor). To measure amplitude quadrature $\hat{X}$ or phase quadrature $\hat{Y}$ of the Pr/Conj fields, the relative phases between them and their LOs are locked by applying the feedback signal from proportional-integral-derivative circuits and high-voltage amplifiers to PZTs.

APPENDIX C: MEASUREMENT OF COVARIANCE MATRIX

To reconstruct the covariance matrix of the CV quantum entangled state, we perform six different measurements on the output optical modes. These measurements include the variances of the amplitude and phase quadratures of Conj and Pr fields $\Delta^2 \hat{X}_{-l,C}$, $\Delta^2 \hat{Y}_{-l,C}$, $\Delta^2 \hat{X}_{l,P}$, $\Delta^2 \hat{Y}_{l,P}$, as well as noise variances of their joint amplitude and phase quadratures $\Delta^2(\hat{X}_{l,P} - \hat{X}_{-l,C})$ and $\Delta^2(\hat{Y}_{l,P} + \hat{Y}_{-l,C})$, respectively. $\Delta^2 \hat{X}_{-l,C}$ and $\Delta^2 \hat{Y}_{-l,C}$ ($\Delta^2 \hat{X}_{l,P}$ and $\Delta^2 \hat{Y}_{l,P}$) are experimentally measured by locking the relative phase of the Conj (Pr) field, and its corresponding local oscillator of BHD₁ (BHD₂) at amplitude quadrature or phase quadrature. $\Delta^2(\hat{X}_{l,P} - \hat{X}_{-l,C})$ and $\Delta^2(\hat{Y}_{l,P} + \hat{Y}_{-l,C})$ are experimentally measured by locking the

relative phases of BHD₁ and BHD₂ at amplitude quadrature or phase quadrature and then subtracting or adding the photocurrents with a radio-frequency subtractor or adder. SNL is achieved by blocking the Pr and/or Conj fields so that only the noise of vacuum is measured. The settings of the SA (Agilent E4411B) are 30-kHz resolution bandwidth, 100-Hz video bandwidth, and zero span at 1.2 MHz.

Figure 6 shows the measured six noise variance levels of the CV entanglement of the OAM multiplexed CV entangled state before being distributed in lossy and noisy channels ($\delta = 0, \eta = 1$). As shown in Fig. 6, the noise levels of amplitude and phase quadratures are $\Delta^2 \hat{X}_{-l,C} = \Delta^2 \hat{Y}_{-l,C} = \Delta^2 \hat{X}_{l,P} = \Delta^2 \hat{Y}_{l,P} = 3.6 \pm 0.1$ dB, and the noise levels of correlated quadratures are $\Delta^2(\hat{X}_{l,P} - \hat{X}_{-l,C}) = \Delta^2(\hat{Y}_{l,P} + \hat{Y}_{-l,C}) = -3.3 \pm 0.1$ dB. Our FWM process works in the amplification regime, so amplitude quadratures and phase quadratures of the Pr/Conj fields show strong correlations and anticorrelations, respectively. It is clear that the degrees of entanglement for CV entangled states carrying topological charges $l = 1$ and $l = 2$ are close to that of their spatially Gaussian counterpart $l = 0$.

With the measured six noise variances, the cross-correlation matrix elements are calculated via

$$\begin{aligned} \text{Cov}(\hat{\xi}_i, \hat{\xi}_j) &= \frac{1}{2} [\Delta^2(\hat{\xi}_i + \hat{\xi}_j) - \Delta^2 \hat{\xi}_i - \Delta^2 \hat{\xi}_j], \\ \text{Cov}(\hat{\xi}_i, \hat{\xi}_j) &= -\frac{1}{2} [\Delta^2(\hat{\xi}_i - \hat{\xi}_j) - \Delta^2 \hat{\xi}_i - \Delta^2 \hat{\xi}_j]. \end{aligned} \quad (\text{C1})$$

In the experiment, we obtain all the covariance matrices of quantum states with different transmission efficiencies and

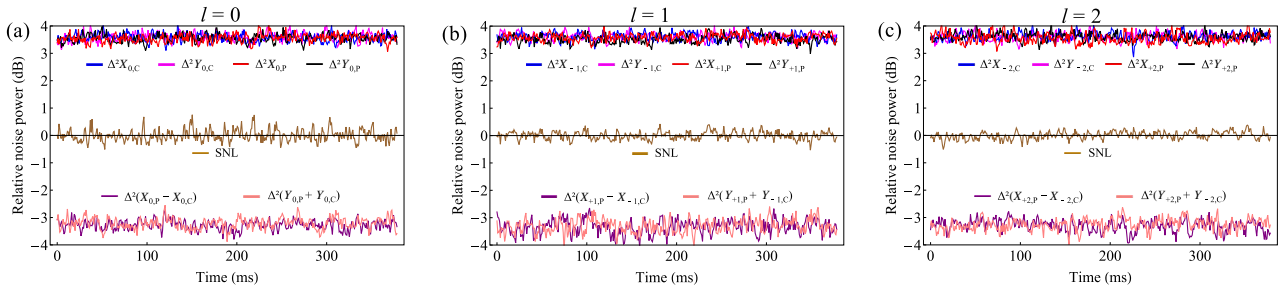


Fig. 6. Measured quantum correlation noises for initially generated OAM multiplexed CV entangled states carrying topological charges (a) $l = 0$, (b) $l = 1$, and (c) $l = 2$, respectively. Brown curve at 0 dB shows the SNL. The other six curves show the noise variances of $\Delta^2 \hat{X}_{-l,C}$, $\Delta^2 \hat{Y}_{-l,C}$, $\Delta^2 \hat{X}_{l,P}$, $\Delta^2 \hat{Y}_{l,P}$, as well as the noise variances of their joint amplitude or phase quadrature $\Delta^2(\hat{X}_{l,P} - \hat{X}_{-l,C})$ and $\Delta^2(\hat{Y}_{l,P} + \hat{Y}_{-l,C})$. These six curves are all normalized to the same SNL. All the measurements are performed at 1.2 MHz. The electronic noise of the BHDs and the background noise from leaked pump fields are subtracted from the SNL and signals, respectively.

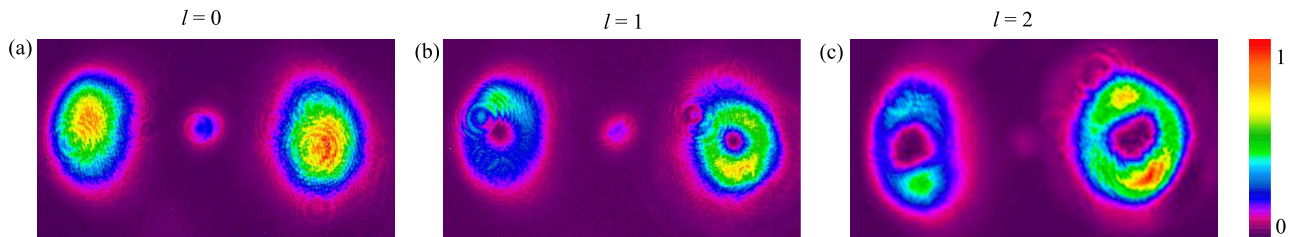


Fig. 7. Images of OAM modes of the Pr beam and Conj beam generated from the FWM process. (a) $l = 0$; (b) $l = 1$; (c) $l = 2$. From left to right: Conj beam, leaked pump beam, and Pr beam.

excess noise and then calculate the smallest symplectic eigenvalue ν of the partially transposed covariance matrices $\mathcal{G}^{A \rightarrow B}(\sigma_{AB})$ and $\mathcal{G}^{B \rightarrow A}(\sigma_{AB})$ to verify whether quantum entanglement and quantum steering exist.

APPENDIX D: SUPPLEMENTAL BEAM PATTERNS

Supplemental beam patterns of the OAM modes of Pr and Conj beams generated from the FWM process are shown in Fig. 7. Pr and Conj beams carrying different topological charges $l = 0$, $l = 1$, and $l = 2$ are generated by removing or switching different VPPs, and their beam patterns are taken after the GT with a camera. It is obvious that the dark hollow patterns of the Laguerre–Gaussian beams are enlarged as the topological charge increases. Furthermore, the beam patterns of Pr and Conj fields are symmetric with respect to the pump beam.

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